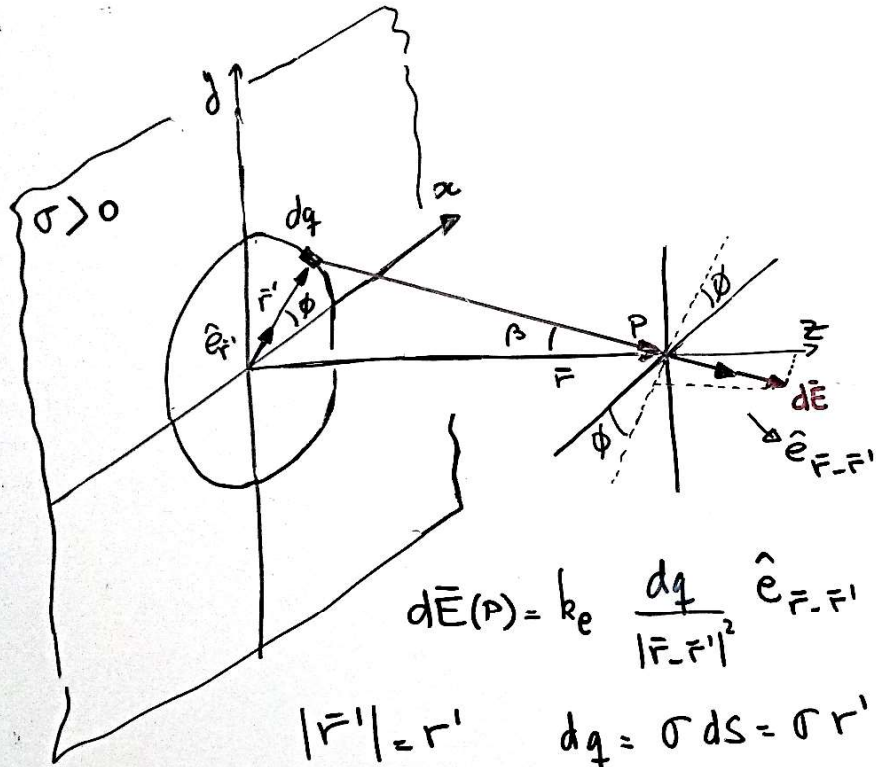


CALCULO DE \vec{E} PARA UN PLANO INFINITO



$$d\vec{E}(P) = k_e \frac{dq}{|\vec{r}-\vec{r}'|^2} \hat{e}_{\vec{r}-\vec{r}'}$$

$$|\vec{r}'| = r' \quad dq = \sigma ds = \sigma r' dr' d\phi$$

$$\vec{r}' = r' \cos \phi \hat{e}_x + r' \sin \phi \hat{e}_y$$

$$\vec{r} = z \hat{e}_z$$

$$\vec{r}-\vec{r}' = z \hat{e}_z - r' \cos \phi \hat{e}_x - r' \sin \phi \hat{e}_y$$

$$|\vec{r}-\vec{r}'|^2 = z^2 + r'^2 \cos^2 \phi + r'^2 \sin^2 \phi = z^2 + r'^2$$

$$\hat{e}_{\vec{r}-\vec{r}'} = \cos \beta \hat{e}_z - \sin \beta \cos \phi \hat{e}_x - \sin \beta \sin \phi \hat{e}_y$$

$$d\vec{E}(P) = k_e \frac{\sigma r' dr' d\phi}{(z^2 + r'^2)} \left[\cos \beta \hat{e}_z - \sin \beta \cos \phi \hat{e}_x - \sin \beta \sin \phi \hat{e}_y \right]$$

Tener en cuenta que,

$$\cos \beta = \frac{z}{(r'^2 + z^2)^{1/2}} \quad \sin \beta = \frac{r'}{(r'^2 + z^2)^{1/2}}$$

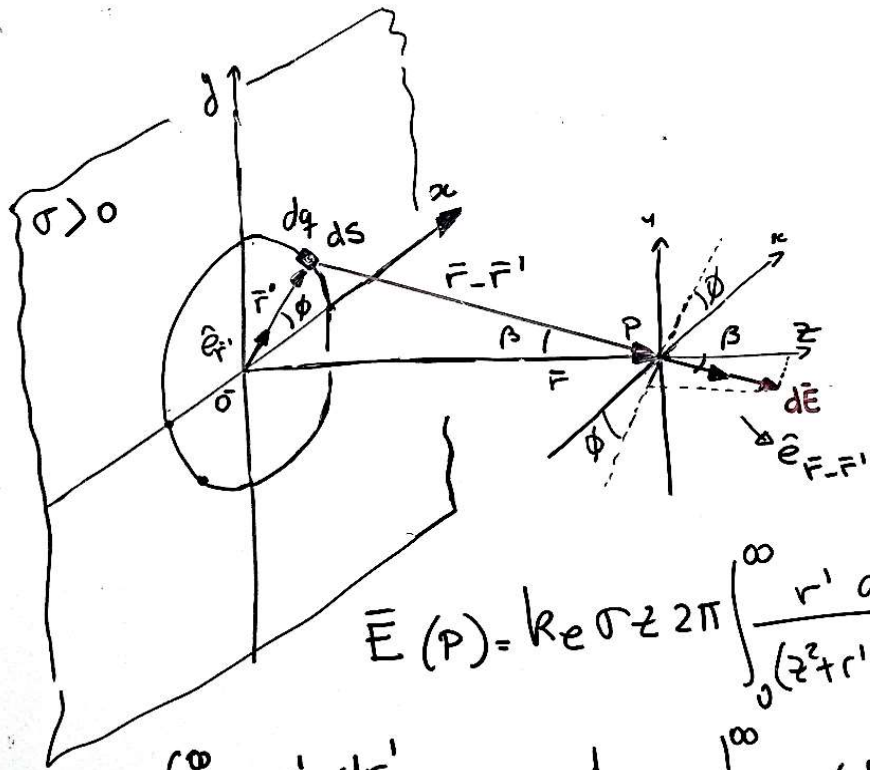
Para encontrar el campo \vec{E} total hay que integrar en todo el plano

$$\vec{E}(P) = k_e \sigma \left[- \int_0^{\infty} \int_0^{2\pi} \frac{r'^2 \cos \phi d\phi dr'}{(r'^2 + z^2)^{3/2}} \hat{e}_x - \int_0^{\infty} \int_0^{2\pi} \frac{r'^2 \sin \phi d\phi dr'}{(r'^2 + z^2)^{3/2}} \hat{e}_y + \int_0^{\infty} \int_0^{2\pi} \frac{z r'}{(r'^2 + z^2)^{3/2}} d\phi dr' \hat{e}_z \right]$$

$$\text{Como } \int_0^{2\pi} \cos \phi d\phi = 0 \quad \text{y} \quad \int_0^{2\pi} \sin \phi d\phi = 0$$

$$\Rightarrow \vec{E}(P) = k_e \sigma z 2\pi \int_0^{\infty} \frac{r'}{(z^2 + r'^2)^{3/2}} dr' \hat{e}_z$$

CALCULO DE \vec{E} PARA UN PLANO INFINITO



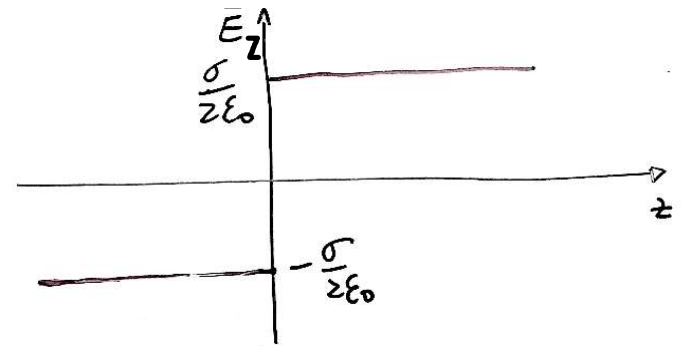
$$\vec{E}(P) = k_e \sigma z 2\pi \int_0^{\infty} \frac{r' dr'}{(z^2 + r'^2)^{3/2}} \hat{e}_z$$

$$\int_0^{\infty} \frac{r' dr'}{(z^2 + r'^2)^{3/2}} = -\frac{1}{(r'^2 + z^2)^{1/2}} \Big|_0^{\infty} = 0 - \left(-\frac{1}{z}\right) = \frac{1}{z}$$

$$k_e = \frac{1}{4\pi\epsilon_0}$$

$$\boxed{\vec{E}(P) = \frac{\sigma}{2\epsilon_0} \frac{z}{z} \hat{e}_z = \frac{\sigma}{2\epsilon_0} \hat{e}_z \quad z > 0}$$

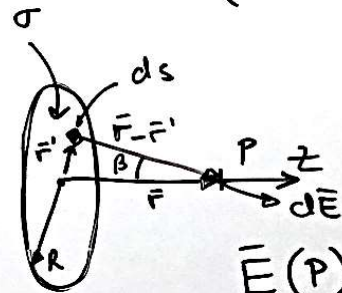
$$\boxed{\vec{E}(P) = -\frac{\sigma}{2\epsilon_0} \hat{e}_z \quad z < 0}$$



CAMPO ELECTRICO DE UN DISCO DE RADIO R

El problema es el mismo que en el caso del plano ∞ , la diferencia son los límites de integración

para r plano $\infty \rightarrow$ disco $0 \leq r' \leq R$

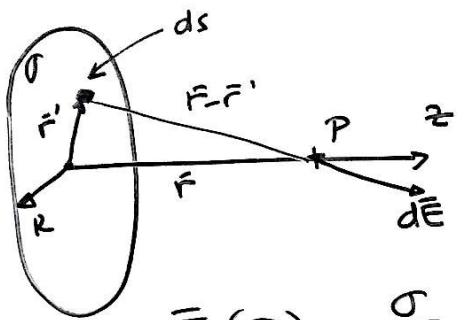


No interesa calcular \vec{E} en el punto P sobre el eje z.

$$\vec{E}(P) = k_e \sigma \int_0^R \int_0^{2\pi} \frac{z r' d\phi dr' \hat{e}_z}{(r'^2 + z^2)^{3/2}}$$

$$= k_e \sigma z 2\pi \left[-\frac{1}{(r'^2 + z^2)^{1/2}} \right]_0^R = k_e \sigma z 2\pi \left[-\frac{1}{(R^2 + z^2)^{1/2}} + \frac{1}{z} \right]$$

$$\boxed{\vec{E}(P) = \frac{\sigma}{2\epsilon_0} \left[1 - \frac{z}{(R^2 + z^2)^{1/2}} \right] \hat{e}_z}$$



$$\vec{E}(P) = \frac{\sigma}{2\epsilon_0} \left[1 - \frac{z}{(R^2 + z^2)^{1/2}} \right] \hat{e}_z$$

Si $z \rightarrow 0$ $\vec{E}(z \rightarrow 0) = \frac{\sigma}{2\epsilon_0} \hat{e}_z$ idem plano ∞

Si $z \rightarrow \infty$ $\vec{E}(z \rightarrow \infty) = \frac{\sigma}{2\epsilon_0} \left[1 - \frac{1}{(1 + (\frac{R}{z})^2)^{1/2}} \right] \hat{e}_z$

$$\vec{E}(z \rightarrow \infty) = \frac{\sigma}{2\epsilon_0} [1 - 1] \hat{e}_z = \vec{0}$$

Si $z \gg R$ $\vec{E}(z) = \frac{\sigma}{2\epsilon_0} \left[1 - \frac{1}{(1 + (\frac{R}{z})^2)^{1/2}} \right] \hat{e}_z$

usamos que $\sqrt{1+x} \approx 1 + \frac{1}{2}x$ si $x \ll 1$

$$\vec{E}(z) = \frac{\sigma}{2\epsilon_0} \left[1 - \frac{1}{1 + \frac{R^2}{2z^2}} \right] \hat{e}_z$$

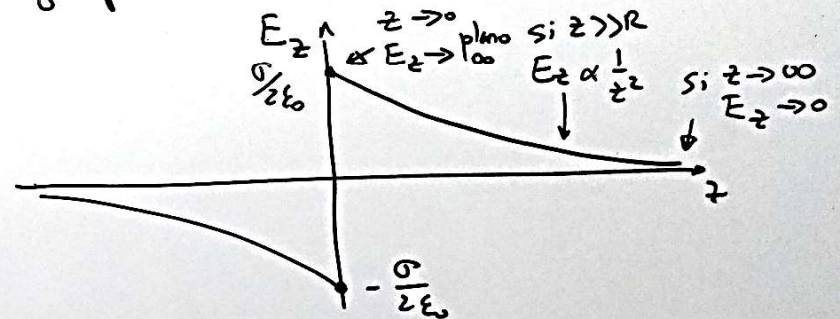
$$= \frac{\sigma}{2\epsilon_0} \left(\frac{1 + \frac{R^2}{2z^2} - 1}{1 + \frac{R^2}{2z^2}} \right) \hat{e}_z = \frac{\sigma}{2\epsilon_0} \frac{R^2/2z^2}{1 + \frac{R^2}{2z^2}} \hat{e}_z$$

$$= \frac{\sigma}{2\epsilon_0} \frac{1}{1 + \frac{2z^2}{R^2}} \hat{e}_z \approx \frac{\sigma}{2\epsilon_0} \frac{R^2}{2z^2} \hat{e}_z$$

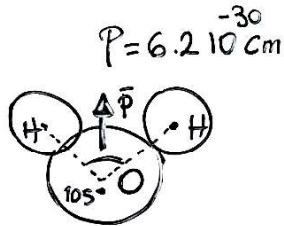
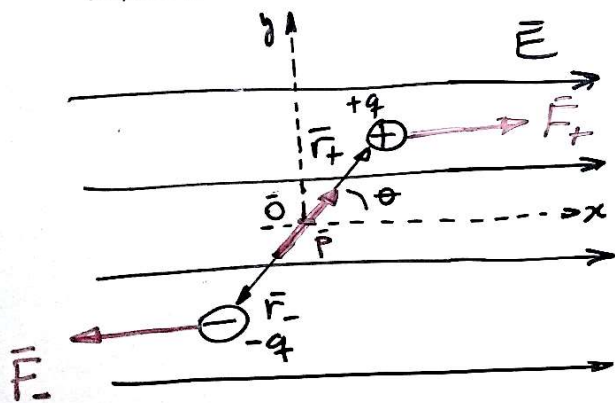
$\frac{2z^2}{R^2} \gg 1$

Como $\sigma = \frac{Q}{\pi R^2} \Rightarrow \vec{E}(z) \approx \frac{Q}{4\pi\epsilon_0} \frac{1}{z^2} \hat{e}_z$

Si graficamos el campo eléctrico



Dipolo en un campo eléctrico \vec{E}



Por un lado tenemos,

$$\vec{p} = p (\cos \theta \hat{e}_x + \sin \theta \hat{e}_y)$$

si $|\vec{r}_+| = |\vec{r}_-| = a$ y $p = 2aq$

$$\vec{p} = 2aq (\cos \theta \hat{e}_x + \sin \theta \hat{e}_y)$$

Por otro lado, la fuerza sobre el dipolo resulta:

$$\vec{F}_+ = q\vec{E} \quad \text{y} \quad \vec{F}_- = -q\vec{E}$$

Por lo que $\vec{F} = \vec{F}_+ + \vec{F}_- = 0$

Sin embargo existe un torque $\vec{\tau}$ respecto al origen de coordenadas \vec{O}

$$\vec{\tau} = \vec{r}_+ \times \vec{F}_+ + \vec{r}_- \times \vec{F}_-$$

donde

$$\vec{r}_+ = a (\cos \theta \hat{e}_x + \sin \theta \hat{e}_y)$$

$$\vec{r}_- = a (-\cos \theta \hat{e}_x - \sin \theta \hat{e}_y)$$

$$\vec{\tau} = a (\cos \theta \hat{e}_x + \sin \theta \hat{e}_y) \times F_+ \hat{e}_x + a (-\cos \theta \hat{e}_x - \sin \theta \hat{e}_y) \times (-F_- \hat{e}_x)$$

$$\vec{\tau} = -a \sin \theta F_+ \hat{e}_z - a \sin \theta F_- \hat{e}_z$$

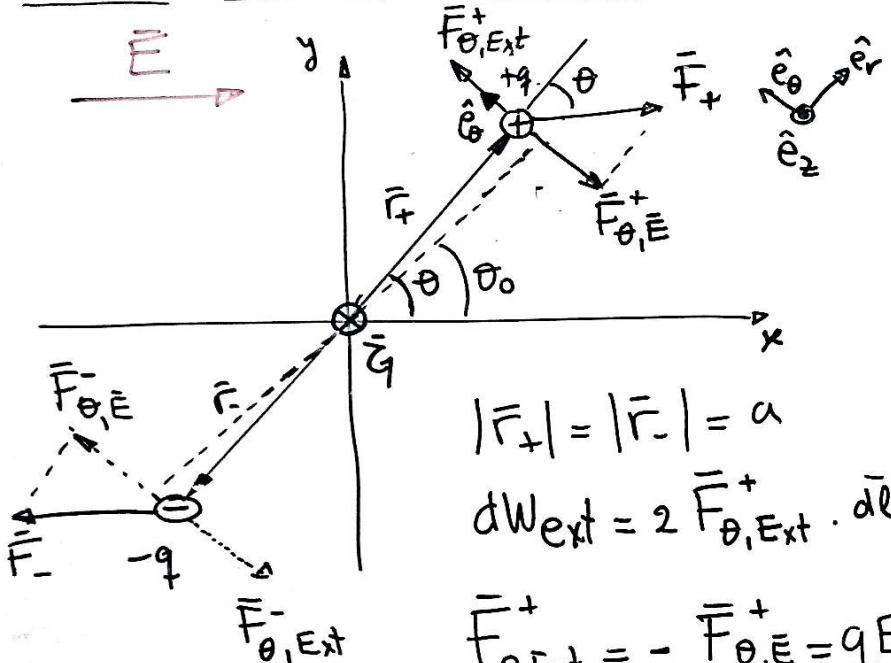
$$F_+ = F_- = F = qE$$

$$\vec{\tau} = -2a \sin \theta F \hat{e}_z$$

$$\vec{\tau} = -2aq \sin \theta E \hat{e}_z = -p E \sin \theta \hat{e}_z$$

$$\boxed{\vec{\tau} = \vec{p} \times \vec{E}}$$

Energía potencial de un dipolo eléctrico



$$|\vec{r}_+| = |\vec{r}_-| = a$$

$$dW_{ext} = 2 \vec{F}_{\theta, Ext}^+ \cdot d\vec{\ell}$$

$$\vec{F}_{\theta, Ext}^+ = -\vec{F}_{\theta, Ext}^- = qE \sin\theta \hat{e}_\theta$$

$$d\vec{\ell} = a d\theta \hat{e}_\theta$$

$$dW_{Ext} = 2 q E \sin\theta a d\theta$$

$$W_{Ext} = 2 q a E \int_{\theta_0}^{\theta} \sin\theta d\theta$$

$$= 2 q a E (-\cos\theta) \Big|_{\theta_0}^{\theta} =$$

$$= -2 q a E (\cos\theta - \cos\theta_0)$$

$$W_{EXT} = \Delta U = U_f - U_0 =$$

$$= -pE (\cos\theta - \cos\theta_0)$$

Si elegimos $\theta_0 = \frac{\pi}{2} \Rightarrow$

$$\Delta U = -pE \cos\theta$$

$$\Delta U = -\vec{p} \cdot \vec{E}$$

Como el sistema será estable en el caso de mínima energía, tenemos ∇U

Si:

- $\theta = 0$
 $\Delta U = -pE$

sistema en equilibrio, mínima energía, estable
- $\theta = \frac{\pi}{2}$
 $\Delta U = 0$

sistema no equilibrado
- $\theta = \pi$
 $\Delta U = pE$

sistema en equilibrio, máxima energía, inestable