



Momento Angular

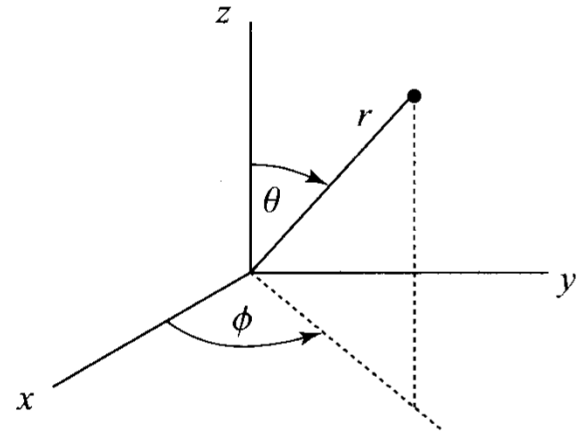
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$$\hat{L}_x = -i\hbar \left(y \frac{\partial}{\partial z} - z \frac{\partial}{\partial y} \right)$$

$$\hat{L}_y = -i\hbar \left(z \frac{\partial}{\partial x} - x \frac{\partial}{\partial z} \right)$$

$$\hat{L}_z = -i\hbar \left(x \frac{\partial}{\partial y} - y \frac{\partial}{\partial x} \right)$$

$$\hat{L}^2 = |\hat{\mathbf{L}}|^2 = \hat{\mathbf{L}} \cdot \hat{\mathbf{L}} = \hat{L}_x^2 + \hat{L}_y^2 + \hat{L}_z^2$$



$$x = r \sin \theta \cos \phi, \quad y = r \sin \theta \sin \phi, \quad z = r \cos \theta$$

$$r^2 = x^2 + y^2 + z^2, \quad \cos \theta = \frac{z}{(x^2 + y^2 + z^2)^{1/2}}, \quad \tan \phi = y/x$$

$$\hat{L}_x = i\hbar \left(\sin \phi \frac{\partial}{\partial \theta} + \cot \theta \cos \phi \frac{\partial}{\partial \phi} \right)$$

$$\hat{L}_y = -i\hbar \left(\cos \phi \frac{\partial}{\partial \theta} - \cot \theta \sin \phi \frac{\partial}{\partial \phi} \right)$$

$$\hat{L}_z = -i\hbar \frac{\partial}{\partial \phi}$$

$$\hat{L}^2 = -\hbar^2 \left(\frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right)$$

Autofunciones Y autovalores (Momento angular)

$$[\hat{L}^2, \hat{L}_z] = 0$$

$$\hat{L}_z Y(\theta, \phi) = bY(\theta, \phi)$$

$$\hat{L}^2 Y(\theta, \phi) = cY(\theta, \phi)$$

$$-i\hbar \frac{\partial}{\partial \phi} Y(\theta, \phi) = bY(\theta, \phi)$$

$$Y(\theta, \phi) = S(\theta)T(\phi)$$

$$-i\hbar \frac{\partial}{\partial \phi} [S(\theta)T(\phi)] = bS(\theta)T(\phi)$$

$$-i\hbar S(\theta) \frac{dT(\phi)}{d\phi} = bS(\theta)T(\phi)$$

$$\frac{dT(\phi)}{T(\phi)} = \frac{ib}{\hbar} d\phi$$

$$T(\phi) = Ae^{ib\phi/\hbar}$$

Es admisible T como función propia?

$$T(\phi) = Ae^{ib\phi/\hbar}$$

$$T(\phi + 2\pi) = T(\phi)$$

$$Ae^{ib\phi/\hbar} e^{ib2\pi/\hbar} = Ae^{ib\phi/\hbar}$$

$$e^{ib2\pi/\hbar} = 1$$

$$e^{i\alpha} = \cos \alpha + i \operatorname{sen} \alpha = 1$$

$$\alpha = 2\pi m$$

$$m = 0, \pm 1, \pm 2, \pm \dots$$

$$2\pi b/\hbar = 2\pi m$$

$$b = m\hbar, \quad m = \dots - 2, -1, 0, 1, \dots$$

$$T(\phi) = Ae^{ib\phi/\hbar}$$

$$T(\phi) = Ae^{im\phi}, \quad m = 0, \pm 1, \pm 2, \dots$$

Normalizando la función encontramos el valor de A

$$\int_0^{2\pi} |T^2| d\phi = 1$$

$$|A| = (2\pi)^{-1/2}$$

$$T(\phi) = \frac{1}{\sqrt{2\pi}} e^{im\phi}, \quad m = 0, \pm 1, \pm 2, \pm \dots$$

$$\hat{L}^2 Y = cY$$

$$Y(\theta, \phi) = S(\theta)T(\phi)$$

$$T(\phi) = \frac{1}{\sqrt{2\pi}} e^{im\phi}, \quad m = 0, \pm 1, \pm 2, \pm \dots$$

$$\hat{L}^2 = -\hbar^2 \left(\frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right)$$

$$-\hbar^2 \left(\frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right) \left(S(\theta) \frac{1}{\sqrt{2\pi}} e^{im\phi} \right) = cS(\theta) \frac{1}{\sqrt{2\pi}} e^{im\phi}$$

$$-\hbar^2 \left(\frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right) \left(S(\theta) \frac{1}{\sqrt{2\pi}} e^{im\phi} \right) = c S(\theta) \frac{1}{\sqrt{2\pi}} e^{im\phi}$$

$$\frac{d^2 S}{d\theta^2} + \cot \theta \frac{dS}{d\theta} - \frac{m^2}{\sin^2 \theta} S = -\frac{c}{\hbar^2} S$$

$$w = \cos \theta$$

$$S(\theta) = G(w)$$

$$(1 - w^2) \frac{d^2 G}{dw^2} - 2w \frac{dG}{dw} + \left[\frac{c}{\hbar^2} - \frac{m^2}{1 - w^2} \right] G(w) = 0$$

$$G(w) = (1 - w^2)^{|m|/2} H(w)$$

$$(1 - w^2) H'' - 2(|m| + 1) w H' + [c \hbar^{-2} - |m| (|m| + 1)] H = 0$$

$$H(w) = \sum_{j=0}^{\infty} a_j w^j$$

Relación de recurrencia:

$$a_{j+2} = \frac{[(j + |m|)(j + |m| + 1) - c/\hbar^2]}{(j + 1)(j + 2)} a_j$$

$$c = \hbar^2(k + |m|)(k + |m| + 1), \quad k = 0, 1, 2, \dots$$

Definimos:

$$l \equiv k + |m|$$

Como

$$|m| \leq l$$

Entonces:

$$m = -l, -l + 1, -l + 2, \dots, -1, 0, 1, \dots, l - 2, l - 1, l$$

Los auto-valores de L^2 son:

$$c = l(l + 1)\hbar^2, \quad l = 0, 1, 2, \dots$$

$$|\mathbf{L}| = [l(l + 1)]^{1/2}\hbar$$

$$S_{l,m}(\theta) = \sin^{|m|}\theta \sum_{\substack{j=1,3,\dots \\ j=0,2,\dots}}^{l-|m|} a_j \cos^j \theta$$

$$P_l^{|m|}(w) \equiv \frac{l}{2^l l!} (1 - w^2)^{|m|/2} \frac{d^{l+|m|}}{dw^{l+|m|}} (w^2 - 1)^l, \quad l = 0, 1, 2, \dots$$

$$P_0^0(w) = 1$$

$$P_1^0(w) = w$$

$$P_1^1(w) = (1 - w^2)^{1/2}$$

$$P_2^0(w) = \frac{1}{2}(3w^2 - 1)$$

$$P_2^1(w) = 3w(1 - w^2)^{1/2}$$

$$P_2^2(w) = 3 - 3w^2$$

$$S_{l,m}(\theta) = \left[\frac{2l + 1}{2} \frac{(l - |m|)!}{(l + |m|)!} \right]^{1/2} P_l^{|m|}(\cos \theta)$$

$$Y_l^m(\theta, \phi) = S_{l,m}(\theta)T(\phi) = \frac{1}{\sqrt{2\pi}} S_{l,m}(\theta) e^{im\phi}$$

$$Y_l^m(\theta, \phi) = \left[\frac{2l+1}{4\pi} \frac{(l-|m|)!}{(l+|m|)!} \right]^{1/2} P_l^{|m|}(\cos\theta) e^{im\phi}$$

$$\hat{L}^2 Y_l^m(\theta, \phi) = l(l+1)\hbar^2 Y_l^m(\theta, \phi), \quad l = 0, 1, 2, \dots$$

$$\hat{L}_z Y_l^m(\theta, \phi) = m\hbar Y_l^m(\theta, \phi), \quad m = -l, -l+1, \dots, l-1, l$$

TABLA 5.1 $S_{l,m}(\theta)$

$l = 0 :$	$S_{0,0} = \frac{1}{2}\sqrt{2}$
$l = 1 :$	$S_{1,0} = \frac{1}{2}\sqrt{6}\cos\theta$ $S_{1,\pm 1} = \frac{1}{2}\sqrt{3}\sin\theta$
$l = 2 :$	$S_{2,0} = \frac{1}{4}\sqrt{10}(3\cos^2\theta - 1)$ $S_{2,\pm 1} = \frac{1}{2}\sqrt{15}\sin\theta\cos\theta$ $S_{2,\pm 2} = \frac{1}{4}\sqrt{15}\sin^2\theta$
$l = 3 :$	$S_{3,0} = \frac{3}{4}\sqrt{14}\left(\frac{5}{3}\cos^3\theta - \cos\theta\right)$ $S_{3,\pm 1} = \frac{1}{8}\sqrt{42}\sin\theta(5\cos^2\theta - 1)$ $S_{3,\pm 2} = \frac{1}{4}\sqrt{105}\sin^2\theta\cos\theta$ $S_{3,\pm 3} = \frac{1}{8}\sqrt{70}\sin^3\theta$

Armónicos esféricos

Table 8.3 The Spherical Harmonics $Y_\ell^{m_\ell}(\theta, \phi)$

$$Y_0^0 = \frac{1}{2\sqrt{\pi}}$$

$$Y_1^0 = \frac{1}{2}\sqrt{\frac{3}{\pi}} \cdot \cos \theta$$

$$Y_1^{\pm 1} = \mp \frac{1}{2}\sqrt{\frac{3}{2\pi}} \cdot \sin \theta \cdot e^{\pm i\phi}$$

$$Y_2^0 = \frac{1}{4}\sqrt{\frac{5}{\pi}} \cdot (3 \cos^2 \theta - 1)$$

$$Y_2^{\pm 1} = \mp \frac{1}{2}\sqrt{\frac{15}{2\pi}} \cdot \sin \theta \cdot \cos \theta \cdot e^{\pm i\phi}$$

$$Y_2^{\pm 2} = \frac{1}{4}\sqrt{\frac{15}{2\pi}} \cdot \sin^2 \theta \cdot e^{\pm 2i\phi}$$

$$Y_3^0 = \frac{1}{4}\sqrt{\frac{7}{\pi}} \cdot (5 \cos^3 \theta - 3 \cos \theta)$$

$$Y_3^{\pm 1} = \mp \frac{1}{8}\sqrt{\frac{21}{\pi}} \cdot \sin \theta \cdot (5 \cos^2 \theta - 1) \cdot e^{\pm i\phi}$$

$$Y_3^{\pm 2} = \frac{1}{4}\sqrt{\frac{105}{2\pi}} \cdot \sin^2 \theta \cdot \cos \theta \cdot e^{\pm 2i\phi}$$

$$Y_3^{\pm 3} = \mp \frac{1}{8}\sqrt{\frac{35}{\pi}} \cdot \sin^3 \theta \cdot e^{\pm 3i\phi}$$

Átomo de Hidrogeno

$$\hat{H} = \hat{T} + \hat{V} = -(\hbar^2/2m)\nabla^2 + V(r)$$

Laplaciano en esfèricas:

$$\nabla^2 = \frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{1}{r^2} \cotg \theta \frac{\partial}{\partial \theta} + \frac{1}{r^2 \text{sen}^2 \theta} \frac{\partial^2}{\partial \phi^2}$$

OP. Momento Angular L^2 en esfèricas:

$$\hat{L}^2 = -\hbar^2 \left(\frac{\partial^2}{\partial \theta^2} + \cot \theta \frac{\partial}{\partial \theta} + \frac{1}{\text{sen}^2 \theta} \frac{\partial^2}{\partial \phi^2} \right)$$

HAMILTONIANO EN ESFERICAS (átomo de Hidrógeno):

$$\hat{H} = -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} \right) + \frac{1}{2mr^2} \hat{L}^2 + V(r)$$

$$\hat{H} = -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} \right) + \frac{1}{2mr^2} \hat{L}^2 + V(r)$$

Se puede demostrar:

$$[\hat{H}, \hat{L}^2] = 0 \quad \text{si } V = V(r)$$

$$[\hat{H}, \hat{L}_z] = 0 \quad \text{si } V = V(r)$$

Entonces:

$$\hat{H}\psi = E\psi$$

$$\hat{L}^2\psi = l(l+1)\hbar^2\psi, \quad l = 0, 1, 2, \dots$$

$$\hat{L}_z\psi = m\hbar\psi, \quad m = -l, -l+1, \dots, l$$

$$-\frac{\hbar^2}{2m} \left(\frac{\partial^2 \psi}{\partial r^2} + \frac{2}{r} \frac{\partial \psi}{\partial r} \right) + \frac{1}{2mr^2} \hat{L}^2 \psi + V(r)\psi = E\psi$$

$$-\frac{\hbar^2}{2m} \left(\frac{\partial^2 \psi}{\partial r^2} + \frac{2}{r} \frac{\partial \psi}{\partial r} \right) + \frac{l(l+1)\hbar^2}{2mr^2} \psi + V(r)\psi = E\psi$$

Proponemos:

$$\psi = R(r)Y_l^m(\theta, \phi)$$

$$-\frac{\hbar^2}{2m} \left(R'' + \frac{2}{r} R' \right) + \frac{l(l+1)\hbar^2}{2mr^2} R + V(r)R = ER(r)$$

Ecuación Radial

Para el caso del átomo de hidrógeno $V(r) = \frac{e^2}{4\pi\epsilon_0 r}$

$$-\frac{\hbar^2}{2m} \left(R'' + \frac{2}{r} R' \right) + \frac{l(l+1)\hbar^2}{2mr^2} R + V(r)R = ER(r)$$

Ecuación Radial para el átomo de hidrógeno

$$\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \frac{2m}{\hbar^2} \left[\left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) - \frac{l(l+1)}{r^2} \right] R = 0$$

$$\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \frac{2m}{\hbar^2} \left[\left(\frac{e^2}{4\pi\epsilon_0 r} + E \right) - \frac{l(l+1)}{r^2} \right] R = 0$$

Las soluciones se llaman **funciones asociadas de LAGUERRE** y las correspondientes energías son

$$E_n = -\frac{m}{2} \left(\frac{e^2}{4\pi\epsilon_0 \hbar} \right)^2 \frac{1}{n^2} = -\frac{E_0}{n^2} \quad n = 1, 2, 3, \dots$$

Misma forma que obtuvo Bohr para el H.

La energía negativa indica que el electrón está ligado al núcleo

$n \rightarrow$ **número cuántico principal**

Y se debe verificar que $l = 0, 1, 2, \dots, (n-1)$

Funciones radiales para el átomo de H

Table 8.4 The Radial Wavefunctions $R_{n\ell}(r)$ of Hydrogen-like Atoms for $n = 1, 2,$ and 3

n	ℓ	$R_{n\ell}(r)$
1	0	$\left(\frac{Z}{a_0}\right)^{3/2} 2e^{-Zr/a_0}$
2	0	$\left(\frac{Z}{2a_0}\right)^{3/2} \left(2 - \frac{Zr}{a_0}\right) e^{-Zr/2a_0}$
2	1	$\left(\frac{Z}{2a_0}\right)^{3/2} \frac{Zr}{\sqrt{3} a_0} e^{-Zr/2a_0}$
3	0	$\left(\frac{Z}{3a_0}\right)^{3/2} 2 \left[1 - \frac{2Zr}{3a_0} + \frac{2}{27} \left(\frac{Zr}{a_0}\right)^2\right] e^{-Zr/3a_0}$
3	1	$\left(\frac{Z}{3a_0}\right)^{3/2} \frac{4\sqrt{2}}{3} \frac{Zr}{a_0} \left(1 - \frac{Zr}{6a_0}\right) e^{-Zr/3a_0}$
3	2	$\left(\frac{Z}{3a_0}\right)^{3/2} \frac{2\sqrt{2}}{27\sqrt{5}} \left(\frac{Zr}{a_0}\right)^2 e^{-Zr/3a_0}$

● Números cuánticos:

● Los tres números cuánticos son:

- n : número cuántico principal
- l : número cuántico orbital
- m_l : número cuántico magnético

● Sus valores permitidos son:

- $n = 1, 2, 3, 4, \dots$
- $l = 0, 1, 2, 3, \dots, n - 1$
- $m_l = -l, -l + 1, \dots, 0, 1, \dots, l - 1, l$

Ejemplo: Indicar la cantidad de estados orbitales que existen para el nivel $n=3$ del átomo de hidrógeno.

$$l = 0, 1, 2, 3, \dots, n-1$$

$$m_l = -l, -l+1, \dots, 0, \dots, l-1, l$$

Para $n=3$ los valores permitidos para l son 0, 1 y 2

Si $l=0$, $m_l=0$

Si $l=1$, hay 3 valores posibles para $m_l = -1, 0, 1$

Si $l=2$, hay 5 valores posibles para $m_l = -2, -1, 0, 1, 2$

l	m_l
0	0
1	-1 0 1
2	-2 -1 0 1 2

Hay 9 estados posibles

Se acostumbra especificar los estados del electrón con distinto momento angular con una letra

■ $l =$	0	1	2	3	4	5 ...
■ Letra =	s	p	d	f	g	h ...

Los estados atómicos usualmente se refieren por sus valores de n y l

Un estado con $n = 2$ y $l = 1$ se llama estado $2p$.

